A Helmholtz energy equation of state for calculating the thermodynamic properties of fluid mixtures

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Abstract

• New Approach

based on highly accurate EOS for the pure components combined at the reduced T, ρ of the mixture

- \blacktriangleright reducing parameters for the T,ρ : composition
- **For simple mixtures**
	- with relatively simple functions
		- very accurate representation
- For nonideal mixtures
	- \triangleright modified reducing func. (T, ρ) +departure func.
		- \rightarrow 0.1% in density, 1% in C_v, 1% in bubble point pressure

Abstract

- **Two applications of mixture model concepts**
	- developed independently in USA and Germany over the same time period
	- including the development of individual equation for each binary system
		- generalization of the model for a wide variety of mixtures
	- Similarities and differences

- \bullet Dimensionless Helmholtz energy, $\alpha = A/RT$
	- Thermodynamic properties of pure fluids with high accuracy over a wide range of T, P
	- All thermodynamic properties from a single mathematical relation
	- A large number of adjustable coeff. (up to 58 terms reported)
	- Reference equation for the calculation of property tables and chart
		- accuracy : restricted only by the accuracy of available measurements

- Due to the complexity Equation with many coefficients • Applied only to pure fluids **High-accuracy EOS**
	- developed for only about 30 pure substances

• Only a few attempts to describe mixture properties using multiparameter EOS

- Plocker et al. (1978)
	- one-fluid theory applied to modified Benedict-Webb-Rubin(mBWR) EOS
	- with a mixing rule for the pseudo-critical temperature
	- ▶ VLE and Enthalpies at high pressures
- Platzer and Maurer (1993)
	- **Exercise 20 Fearth Bender equation**
	- **Pericular Franchista Franchista Exercise** Properties
- Huber and Ely (1994)
	- riangleright extended corresponding states model based on mBWR EOS

 \bullet Mixture model (pure \rightarrow mixture)

based on high-accuracy pure fluid EOS for the pure components

Pressure-explicit EOS

• Simple integration of pressure

generally explicit in Helmholtz energy

 \blacktriangleright adopted for the mixture

• combination rules for the Helmholtz energy function of mixture components (not mixing rules)

• not required a fixed structure of the pure fluid equation

single mathematical expression

• General form of Helmholtz energy model

- independently developed by both authors during overlapping periods of time
- Tillner-Roth (1993)
	- focused on developing an accurate formulation for individual binary mixtures for a large amount of accurate experimental data
- **Lemmon** (1996)
	- focused on developing a generalized model capable of accurate property calculation for a large number of fluid mixtures

• Purpose

- additional detail and background for the models
	- especially ideal mixture
- difference between the models

Helmholtz energy

• EOS for the pure fluids \triangleright ideal part + residual part = $\frac{A}{RT} = \alpha(\tau, \delta) \, \alpha^{\mathrm{o}}(\tau, \delta) + \alpha^{\mathrm{r}}(\tau, \delta)$ 21*kk*

$$
\alpha^{0}(\tau,\delta) = \ln \delta + c_1 \ln \tau + \sum_{k=2}^{k} c_k \tau^{m_k} + \sum_{k=k}^{k} c_k \ln[1 - \exp(-a_k \tau)]
$$

$$
\alpha^{\mathrm{r}}(\tau,\delta) = \sum_{k=k}^{k} c_k \delta^{n_k} \tau^{m_k} + \sum_{k=k}^{k} c_k \delta^{n_k} \tau^{m_k} \exp(-\delta^{e_k})
$$

$$
\tau = \frac{T_n}{T}, \delta = \frac{\rho}{\rho_n} = \frac{V_n}{V}
$$

 \blacktriangleright reducing parameter T_n , V_n : critical properties

Helmholtz energy

• Transforming into the Helmholtz energy form by integrating the relation

$$
p = \rho RT \left[1 + \delta \left(\frac{\partial \alpha^{r}}{\partial \delta} \right)_{\tau} \right]
$$

 \bullet α of mixture : composition dependence

$$
\alpha = \alpha^0 + \alpha^r
$$

• Ideal gas mixture : analytically from functions for the ideal gas properties of the pure fluids

$$
\alpha^0 = \alpha^0(T, V, \mathbf{x}) = \sum_{i=1}^l x_i \alpha_i^0(T, V) + \sum_{i=1}^l x_i \ln x_i
$$

• consistent with ideal parts of the pure fluid equation with independent reduced variable τ , δ

• Residual Helmholtz energy

$$
\alpha^{r}(\tau,\delta,\mathbf{x}) = \sum_{i=1}^{l} x_{i} \alpha_{i}^{r}(\tau,\delta) + \Delta \alpha^{r}(\tau,\delta,\mathbf{x})
$$

linear combination of the pure fluid residual energy \blacktriangleright departure function, $\Delta \alpha^r$ \rightarrow interaction of different species in the mixture ▶ Corresponding states principle shape factors : ratios of the critical properties

$$
T_r = T \cdot h = T \frac{T_{n,r}}{T_n}, \quad \tau_r = \frac{T_{n,r}}{T_r} = \frac{T_n}{T} = \tau
$$

Table 1

• Reducing parameter of the pure fluid equation \neq reducing functions of the mixture *pseudo-critical* parameter τ, δ **Examposition dependence**

$$
\tau(\mathbf{x}) = \frac{T_n(\mathbf{x})}{T}, \delta(\mathbf{x}) = \frac{V_n(\mathbf{x})}{V}
$$

 \rightarrow as x_i approaches 1

$$
T_n(\mathbf{x}) \to T_{n,i}, V_n(\mathbf{x}) \to V_{n,i}
$$

or

$$
\tau(\mathbf{x}) \to \tau_i, \delta(\mathbf{x}) \to \delta_i
$$

Linear mixture model

• Linear mixture model

 \blacktriangleright linear combinations of the pure fluid parameters

$$
T_n(\mathbf{x}) = \sum_{i=1}^l x_i T_{n,i}, \quad V_n(\mathbf{x}) = \sum_{i=1}^l x_i V_{n,i}
$$

 \blacktriangleright departure function $\Delta \alpha^r$ omitted no adjustable parameters

• 1,1,1,2-tetrafluoroethane(R-134a) +difluoromethane(R-32) **Example 2 Sepandical Solution** ▶ VLE : excellent results (Fig.1a) Density in subcritical range (Fig.2a) • vapor, liquid density : good agreement Density in supercritical range (Fig.2b) **Plarger deviation Isochoric heat capacity**

 $\rightarrow \pm 1\%$ (about the experimental uncertainty)

● water+ammonia

• mixture of polar components showing large mixing effects, wide two-phase region

 \blacktriangleright VLE (Fig.1b)

dew curve : well

 \star high conc. in the vapor phase \rightarrow ideal at lower pressures

bubble curve : large deviation

- Density (Fig.2c)
	- \rightarrow -1% \sim +3% deviation
- Excess enthalpy (Fig.3c)

→ poor representation : related to VLE

● ethane+carbon dioxide azeotropic behavior \blacktriangleright VLE (Fig.1c) • incorrect prediction ▶ Isochoric heat capacity (Fig.3b) **less accurate** \rightarrow >2% deviation outside of the critical region

Fig. 1. Experimental and calculated VLE state points from the linear mixture model. (a) R-134a + R-32: \circ , \bullet Holcomb et al. [6]. (b) Water + ammonia: \bigcirc , \bullet Smolen et al. [26]. (c) Ethane + carbon dioxide: \bigcirc , \bullet Fredenslund and Mollerup [3].

Fig. 2. Density deviations between experimental data and densities calculated from the linear mixture model. (a) Subcritical R-134a + R-32 ($T = 333.15$ K, $x ≈ 0.5$): \times , \circ Kleemiss [9]; ■ Magee and Haynes [15]. (b) Supercritical R-134a + R-32 $(363 < T$ (K) < 413, $x \approx 0.5$): \times , \circ Kleemiss [9]. (c) Liquid water+ammonia (273 < T (K) < 333, $p < 40$ MPa): O Harms-Watzenberg [5].

Fig. 3. Deviations between measured caloric properties and values calculated from the linear mixture model. (a) Liquid R-134a + R-32 C_V ($x \approx 0.5$): \Box Magee and Haynes [15]. (b) Isochoric heat capacity for ethane + carbon dioxide: \Diamond Magee [13]. (c) Excess enthalpy in liquid water + ammonia ($T = 298.15$ K, $p = 5$ MPa): \bullet Staudt [28].

• Modification of the reducing functions

$$
T_n(\mathbf{x}) = \sum_{i=1}^{l} \sum_{j=1}^{l} x_i x_j T_{n,ij} \qquad T_{n,ij} = k_{T,ij} \frac{1}{2} (T_{n,ii} + T_{n,jj})
$$

$$
V_n(\mathbf{x}) = \sum_{i=1}^{l} \sum_{j=1}^{l} x_i x_j V_{n,ij} \qquad V_{n,ij} = k_{V,ij} \frac{1}{2} (V_{n,ii} + V_{n,jj})
$$

 $\blacktriangleright k_{T,ij}$, $k_{V,ij}$: adjustable parameter \blacktriangleright Effect of varying $k_{T,ii}$, $k_{V,ii}$ \rightarrow VLE of ethane+carbon dioxide (Fig.4) ($k_{T,ij}$, $k_{V,ij}$) $-C_v$: 1.4% deviation (Fig.5)

Fig. 4. Influence of the parameters $k_{T,12}$ and $k_{F,12}$ in the quadratic reducing function on the VLE of ethane+carbon dioxide: O, \bullet Fredenslund and Mollerup [3]; solid line: $k_T = 1.0$, $k_V = 1.0$.

Fig. 5. Deviations between ethane+carbon dioxide isochoric heat capacities and values calculated from the Helmholtz energy model with $k_{T,12} = 0.9$: \diamond Magee [13].

• Modification by Tilner-Roth \triangleright supplement the quadratic expression with β_{τ} , β_{ν}

$$
T_n(\mathbf{x}) = x_1^2 T_{n,1} + x_2^2 T_{n,2} + 2x_2 (1 - x_2^{\beta_T}) T_{n,12}
$$

$$
V_n(\mathbf{x}) = x_1^2 V_{n,1} + x_2^2 V_{n,2} + 2x_2 (1 - x_2^{\beta_V}) V_{n,12}
$$

 \triangleright effect of exponents on the shape (Fig.6) influence of $T_{n,12}$ or $V_{n,12}$ • asymmetric with respect to composition • Additional flexibility for tuning of the mixture model

Fig. 6. Influence of the exponent β_T on the function $x(1 - x^{\beta_T})$.

• Modification by Lemmon

• modified the linear reducing function

$$
T_{n}(\mathbf{x}) = \sum_{i=1}^{l} x_{i} T_{n,i} + \sum_{i=1}^{l-1} \sum_{j=i+1}^{l} x_{i}^{\beta_{ij}} x_{j} \zeta_{ij}
$$

$$
V_{n}(\mathbf{x}) = \sum_{i=1}^{l} x_{i} V_{n,i} + \sum_{i=1}^{l-1} \sum_{j=i+1}^{l} x_{i} x_{j} \xi_{ij}
$$

 $\sum_i \xi_{ii}$, ξ_{ii} , β_{ii} : modify the shape of the reducing parameter $\beta_{ij} = 1, \quad \zeta_{ij} = (T_{n,i} + T_{n,j})(k_{T,ij} - 1)$ $\xi_{ij} = (V_{n,i} + V_{n,j})(k_{V,ij} - 1)$ $V \cdot + V \cdot \mathcal{N}$ $T \t+T \t\t k$

 \rightarrow for ethane+CO₂, $\zeta_{ii} = -67.4 \text{K}$ ($k_{T,ii} = 0.89$)

Departure function

• Departure function

$$
\Delta \alpha^{r}(\tau, \delta, \mathbf{x}) = \sum_{i=1}^{l-1} \sum_{j=i+1}^{l} x_{i} x_{j} \Delta \alpha_{ij}(\tau, \delta, x_{i}, x_{j})
$$

- only to mixture properties, no effect at the pure fluid **limits**
- magnitude of departure function
	- \rightarrow generally one order smaller than the residual α ^r
- modeling : regression analysis by Wagner
	- → a set of of experimental data → lowest standard deviation→bank of terms→subset of term for Δα*ij*
	- optimization : modified Marquardt-Fletcher algorithm

 \bullet Introduction an exponent γ in the departure function for a binary system

$$
\Delta \alpha^{r}(\tau, \delta, x) = x_2(1 - x_2^{\gamma})\Delta \alpha_{ij}(\tau, \delta, x_i, x_j)
$$

asymmetric influence

→ very effective for simple system, R-152+R134a

reduce the number of terms in the departure function

• applied only to binary systems to date

▶ For water+ammonia,

$$
x\delta^{n_k}\tau^{m_k}\exp\left(-\delta^{e_k}\right)
$$
 and $x^2\delta^{n_k}\tau^{m_k}\exp\left(-\delta^{e_k}\right)$

• Results for R-134a+R-32

Five-term departure function

$$
\Delta \alpha^r(\tau, \delta, \mathbf{x}) = x(1 - x^{\gamma}) \left[\sum_{k=1}^2 c_k \delta^{n_k} \tau^{m_k} + \sum_{k=3}^5 c_k \delta^{n_k} \tau^{m_k} \exp(-\delta^{e_k}) \right]
$$

 \blacktriangleright Fig.7 (cf. Fig.2b)

 \rightarrow density at liquid & vapor region : $\pm 0.1\%$

 \rightarrow density at critical region : -6 ~ \pm 0.4%

- Results for water+ammonia
	- ▶ 14-term departure function : 19 adjustable parameters

$$
\frac{\Delta \alpha^r(\tau, \delta, \mathbf{x})}{x(1-x^{\gamma})} = c_1 \delta^{n_1} \tau^{m_1} + \sum_{k=2}^6 c_k \delta^{n_k} \tau^{m_k} \exp\left(-\delta^{e_k}\right) + x \sum_{k=7}^{13} c_k \delta^{n_k} \tau^{m_k} \exp\left(-\delta^{e_k}\right) + c_{14} x^2 \delta^{n_{14}} \tau^{m_{14}} \exp\left(-\delta^{e_{14}}\right)
$$

 \blacktriangleright Fig.8 (cf. Fig.1b, Fig.3c) improved representation

Fig. 8. Vapor-liquid equilibrium and excess enthalpy of water+ammonia. (a) VLE data: O Rizvi and Heidemann [21]; ○ Sassen et al. [24]. (b) Excess enthalpies (298 < T (K) < 373, $5 < p$ (MPa) < 12): ● Staudt [28]; solid line: Helmholtz energy model including departure function; dashed line: linear mixture model.

• Generalized departure function

$$
\Delta \alpha^{\mathrm{r}}(\tau, \delta, \mathbf{x}) = \sum_{i=1}^{l-1} \sum_{j=i+1}^{l} x_i x_j F_{ij} \Delta \alpha_{ij}(\tau, \delta)
$$

$$
\Delta \alpha_{ij}(\tau, \delta) = \sum_{k=1}^{l} c_k \delta^{n_k} \tau^{m_k}
$$

$$
\blacktriangleright
$$
 c_k , n_k , m_k : fixed the same for all mixture

=

k

1

- \blacktriangleright *F_{ij}*, ζ _{*ij*}, ζ _{*ij*}, β _{*ij*} : adjustable parameters from (p, V, T, x) \rightarrow 0.2% in density, 1% in heat capacities
	- 1-2% bubble point pressure at critical temp.
	- **→**5-10% bubble point pressure at critical temp. father apart

• Results

- propane $+ n$ -butane : Fig.10
	- **Example 1.04%** average absolute deviation in density
	- 0.4% average absolute deviation in bubble point pressure
- **Independent Uncertainty of the equation reported**
	- \rightarrow 0.2% in density
	- \rightarrow 1% in heat capacities
- \triangleright Fig.11 : ternary refrigerant mixture

Fig. 10. The mixture propane + n -butane. (a) Deviations between liquid densities and values predicted from the Helmholtz energy model: \bullet Parrish [16] (283 < T (K) < 333, 0.15 < x_{propane} < 0.75); \circ Magee [14] (244 < T (K) < 400, $x_{\text{propane}} \approx 0.6$). (b) Deviations between bubble point pressures and values predicted from the Helmholtz energy model: \bullet Kay [8] (332 < T $(K) < 425$).

Fig. 11. Deviations between experimental densities for the ternary system $R-134a + R-125 + R-32$ and values predicted from the Helmholtz energy model: \bullet Kleemiss [9] (243 < T (K) < 353); \circ Magee and Haynes [15] (200 < T (K) < 400).

Conclusions

- EOS for mixture in terms of the Helmholtz energy has been developed and applied to several fluid systems by two different research groups
- By Tillner-Roth, model of single binary systems for which comprehensive accurate data are available
- By Lemmon, generalized model was presented which accurately predicts thermodynamic properties of mixtures
- Further research for more nonideal systems and complex multicomponent mixtures