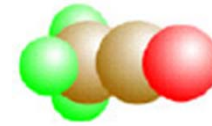


Applied Statistical Mechanics
Lecture Note - 12

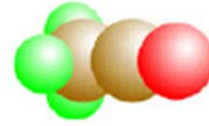


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Advanced Topics in Molecular Monte Carlo Method

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강정원

Contents



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- I. Free Energy Calculations
- II. Gibbs Ensemble Method
- III. Biased Sampling Method

Introduction



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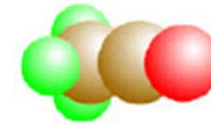
- Previous simulation methods : Properties of a single isolated phase
- Molecular Simulation Techniques can be extended to multiple phases
 - Thermodynamic Phase Coexistence : Two or more phases are equally stable
 - Condition of Phase Coexistence

$$T^{\alpha} = T^{\beta}$$

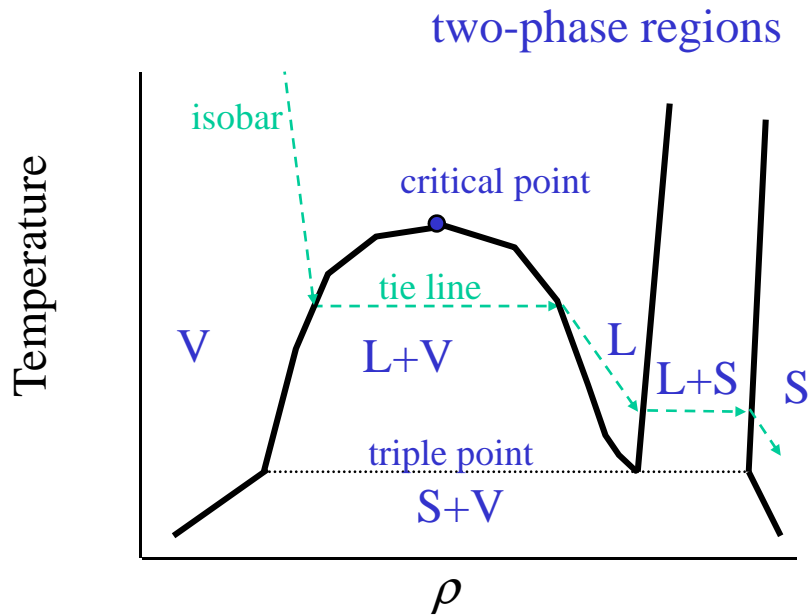
$$p^{\alpha} = p^{\beta}$$

$$\mu_i^{\alpha} = \mu_i^{\beta} \quad i = 1, \dots, C$$

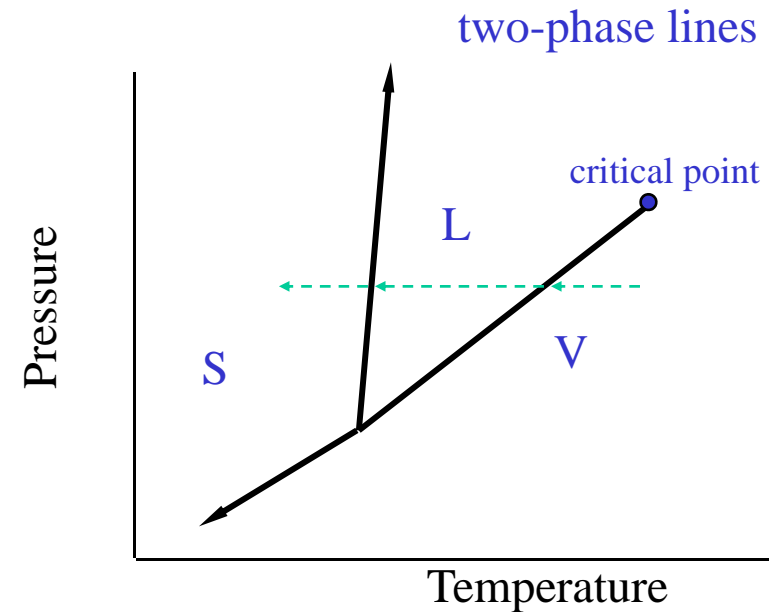
Phase Diagrams



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ρ -T diagram



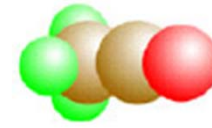
TP diagram

■ Gibbs Phase Rule

$$\square F = 2 + C - P$$

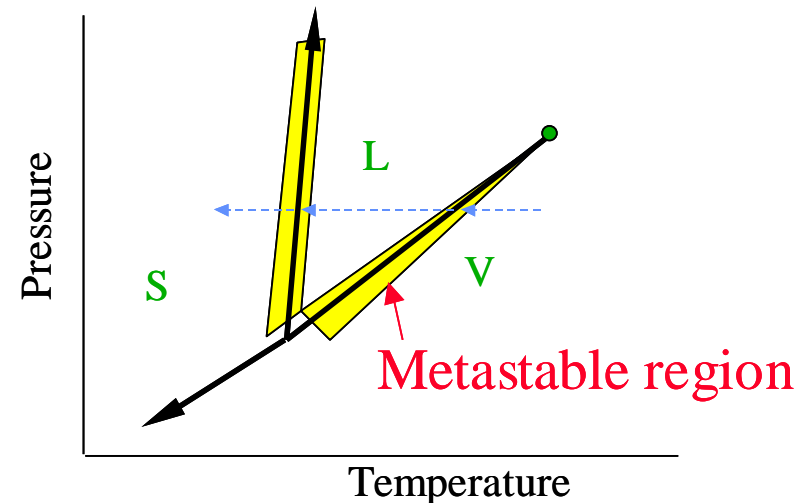
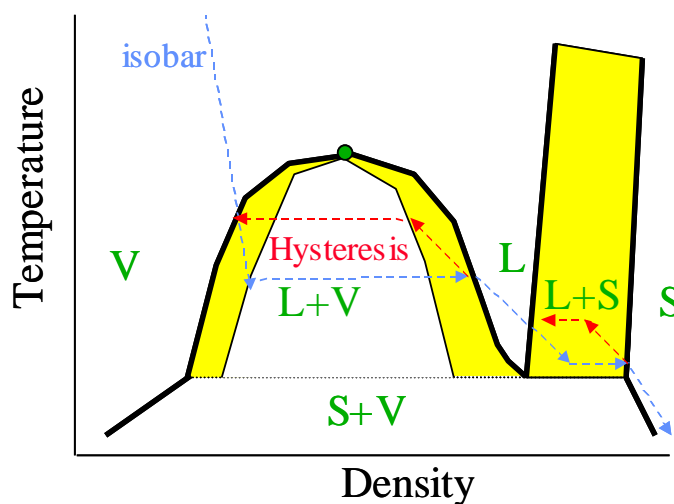
- F : Degree of Freedom
- C : Number of Components
- P : Number of Phases

Straightforward Simulation

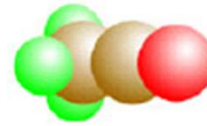


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- Change T or P for given model system and wait for a phase transformation occur
 - Major Drawbacks
 - Hysteresis : Irreversible phase transformation
 - Large free energy barrier at interface \rightarrow Depends on the size of the interface \rightarrow Depends on the choice of simulation system



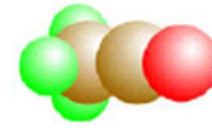
Alternative Methods



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- Gibbs Ensemble Method (Panagiotopoulos, 1987)
- Gibbs – Duhem Integration Method (Kofke, 1993)

Why free energies are important in Phase Coexistence ?



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■ From 2nd Law of Thermodynamics

□ At equilibrium, S is maximum for (N, V, E)

□ (N, V, E) System : A is minimum

$$A = U - TS \text{ (Helmholtz Free Energy)}$$

□ (N, P, T) System : G is Minimum

$$G = H - TS \text{ (Gibbs Free Energy)}$$

■ Equilibrium Condition

$$A^\alpha = A^\beta$$

$$G^\alpha = G^\beta$$

Free energy cannot be directly measured in simulation



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■ Helmholtz free energy

$$A = -kT \ln Q(N, V, T) = -kT \left(\frac{\int dp^N dr^N \exp(-\beta H(p^N, r^N))}{\Lambda^{dN} N!} \right)$$

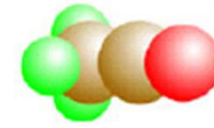
- Not the form of canonical average over phase space
- Depends directly on the available volume in the phase space
- Cannot be directly measured in real experiment, too.

■ Derivatives of free energy

$$\left(\frac{\partial A}{\partial V} \right)_{N, T} = -p \qquad \left(\frac{\partial A/T}{\partial 1/T} \right)_{V, T} = -E$$

→ Find a reversible path in V,T plane and perform integration

Kirkwood's coupling parameter method



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- Assume U depends on the coupling parameter λ :

$$U(\lambda) = (1-\lambda)U_I + \lambda U_{II} = U_I + \lambda(U_{II} - U_I)$$

- Partition function for potential energy function

$$Q(N, V, T, \lambda) = \frac{1}{\Lambda^{3N} N!} \int dr^N \exp(-\beta U(\lambda))$$

$$\begin{aligned} \left(\frac{\partial A(\lambda)}{\partial \lambda} \right)_{N, V, T} &= -\frac{1}{\beta} \frac{\partial}{\partial \lambda} \ln Q(N, V, T, \lambda) = -\frac{1}{\beta Q(N, V, T, \lambda)} \frac{\partial Q(N, V, T, \lambda)}{\partial \lambda} \\ &= \frac{\int dr^N \left(\frac{\partial U(\lambda)}{\partial \lambda} \right) \exp(-\beta U(\lambda))}{\int dr^N \exp(-\beta U(\lambda))} = \left\langle \frac{\partial U(\lambda)}{\partial \lambda} \right\rangle_{\lambda} \longrightarrow \text{Ensemble average of} \\ &\hspace{15em} \text{with potential function } U(\lambda) \end{aligned}$$

Kirkwood's coupling parameter method



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■ Free energy difference

$$A(\lambda = 1) - A(\lambda = 0) = \int_{\lambda=0}^{\lambda=1} d\lambda \left\langle \frac{\partial U(\lambda)}{\partial \lambda} \right\rangle$$

■ All free energy methods are based on calculation of free energy differences

■ Example :

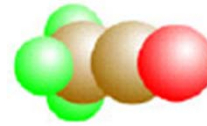
□ Volume of R can be measured as a fraction of total volume

- Sample the reference system
- keep an average of the fraction of time occupying the target system

$$\frac{\Omega_R}{\Omega_\Gamma} = \langle s(\Gamma) \rangle_\Gamma$$

$$S_R - S_\Gamma = k \ln(\Omega_R / \Omega_\Gamma)$$

Chemical Potentials



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■ Chemical Potentials

$$\mu_a = \left(\frac{\partial G}{\partial N_a} \right)_{T,P,N_{b \neq a}} = \left(\frac{\partial A}{\partial N_a} \right)_{T,V,N_{b \neq a}} = -T \left(\frac{\partial S}{\partial N_a} \right)_{V,E,N_{b \neq a}}$$

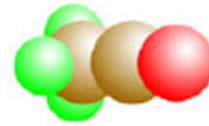
$$Q(N, V, T) = \frac{V^N}{\Lambda^{dN} N!} \int_0^1 \dots \int_0^1 ds^N \exp(-\beta U(s^N; L))$$

$$s^N = r^N / L$$

$$A(N, V, T) = -kT \ln Q = -kT \ln \left[\frac{V^N}{\Lambda^{dN} N!} \right] - kT \ln \left\{ \int_0^1 \dots \int_0^1 ds^N \exp(-\beta U(s^N; L)) \right\}$$

$$= A_{id}(N, V, T) + A_{ex}(N, V, T)$$

Chemical Potentials



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- For sufficiently large N,

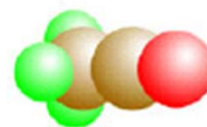
$$\mu = -kT \ln(Q_{N+1} / Q_N)$$

$$\mu = -kT \ln(Q_{N+1} / Q_N) = -kT \ln \left[\frac{V / \Lambda^d}{N+1} \right] - kT \ln \left\{ \frac{\int_0^1 \dots \int_0^1 ds^{N+1} \exp(-\beta U(s^{N+1}))}{\int_0^1 \dots \int_0^1 ds^N \exp(-\beta U(s^N))} \right\}$$

$$= \mu_{id}(\rho) + \mu_{ex}$$

$$\mu_{ex} = -kT \ln \int ds^{N+1} \langle \exp(-\beta \Delta U) \rangle_N$$

Widom's Test Particle Insertion Method



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- Widom (1963)
- A “ghost particle” is randomly inserted into the ensemble and calculating the energy of its interaction
- The test particle is a “ghost”, it does not affect the properties of real molecule
- In principle , this method can be used in any simulation system for the calculation of chemical potential

Implementation of Widom's Method



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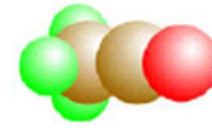
- Carry out conventional NVT or NPT Monte Carlo Simulation of N particles
- At frequent interval during simulation, randomly generate a coordinate, s_{N+1} uniformly over unit cube
- For given s_{N+1} , compute :

$$\mu_{ex} = -kT \ln \langle \exp(-\beta \Delta U) \rangle \longrightarrow \text{NVT ensemble}$$

$$\mu_{ex} = -kT \ln \left[\frac{\langle V \exp(-\beta \Delta U) \rangle}{\langle V \rangle} \right] \longrightarrow \text{NPT ensemble}$$

- Average \rightarrow chemical potential

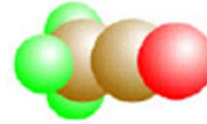
Note on Widom's Method



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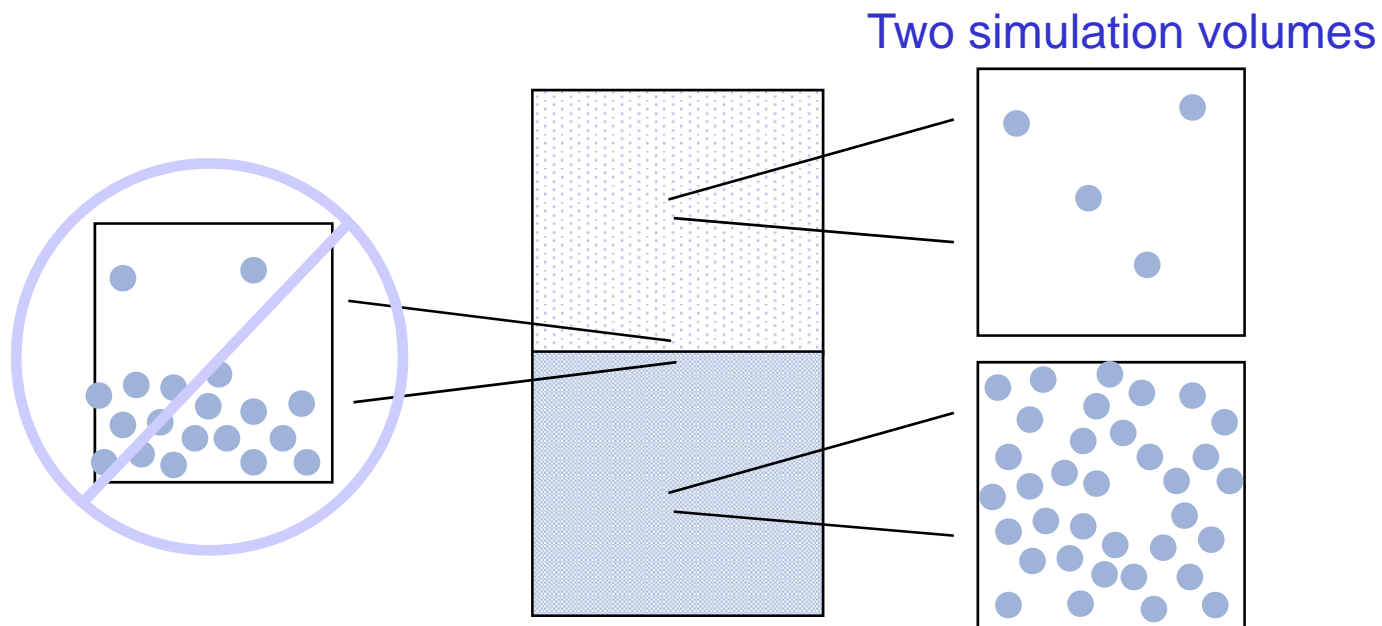
- Widom's method can be used as a verification that equilibrium has been attained
- At high density, some difficulties are encountered
 - Difficult to insert a particle at given location

Gibbs Ensemble Method

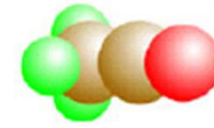


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- Proposed by Panagiotopoulos (1987,1988)
- Simulation method without interface
- Thermodynamic contact without physical contact



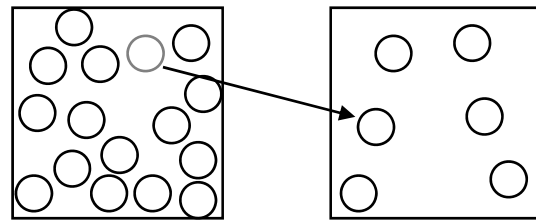
Gibbs Ensemble Method



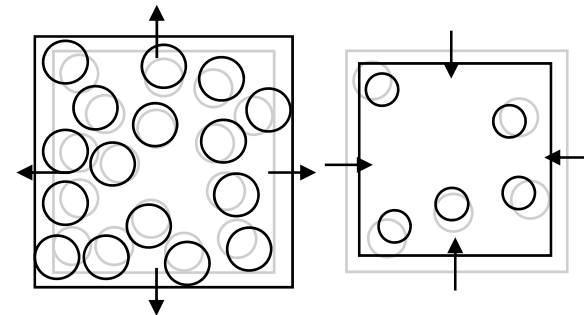
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- MC simulation includes moves that couples two simulation volumes

Particle exchange equilibrates chemical potential

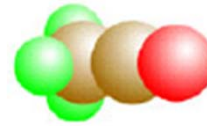


Volume exchange equilibrates pressure



Incidentally, the coupled moves enforce mass and volume balance

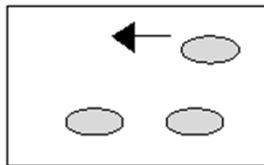
Gibbs Ensemble Method



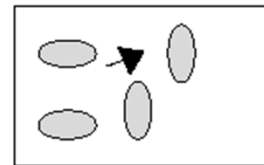
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(1) Attempt molecular displacement

Box 1: N_1, V_1, T

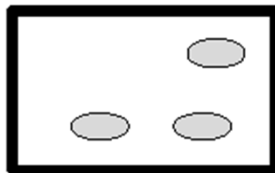


Box 2: N_2, V_2, T

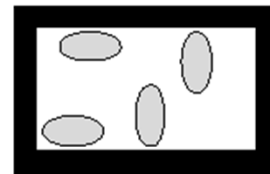


(2) Attempt volume fluctuation

Box 1: $N_1, V_1 + \Delta V, T$

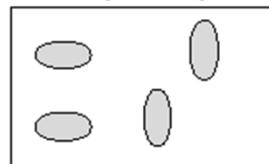


Box 2: $N_2, V_2 - \Delta V, T$

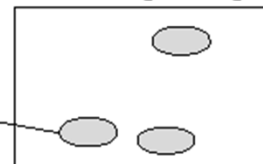


(3) Attempt molecular transfer

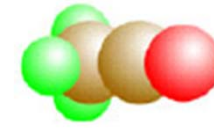
Box 1: $N_1 + 1, V_1, T$



Box 2: $N_2 - 1, V_2, T$



Gibbs Ensemble Method



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- Acceptance of three moves → Governed by Pseudo Boltzman Factor

- Displacement movement

$$\Delta Y_{disp} = \Delta E_{\alpha} + \Delta E_{\beta}$$

- Volume change

$$\Delta Y_{disp} = \Delta E_{\alpha} + \Delta E_{\beta} - N_{\alpha} kT \ln \frac{V_{\alpha} + \Delta V_{\alpha}}{V_{\alpha}} - N_{\beta} kT \ln \frac{V_{\beta} + \Delta V_{\beta}}{V_{\beta}} + P(\Delta V_{\alpha} + \Delta V_{\beta})$$

- Molecular transfer

$$\Delta Y_{disp} = \Delta E_{\alpha} + \Delta E_{\beta} - kT \ln \frac{V_{\beta}(N_{\alpha} + 1)}{V_{\alpha} N_{\beta}}$$

Gibbs Ensemble - Algorithm



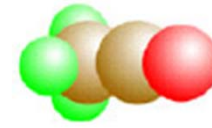
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```
Part 1      Initialisation:
            Specify the number of cycles ( $nCycles$ ), number of particles in
            the boxes ( $nBox1$ ,  $nBox2$ ), number of volume ( $nVol$ ), transfer
            attempts ( $nTrans$ ), and the total number of moves
            ( $nTotal = nBox1 + nBox2 + nVol + nTrans$ ).

Part 2      Generate Markov chain:
            loop  $i \leftarrow 1 \dots nCycles$ 
                loop  $j \leftarrow 1 \dots nTotal$ 
                    if ( $rand() \leq nBox1/nTotal$ ) //Decide which move to perform
                        Displace particle in box 1.
                    else if ( $rand() \leq (nBox1 + nBox2)/nTotal$ )
                        Displace particle in box 2.
                    else if ( $rand() \leq (nBox1 + nBox2 + nVol)/nTotal$ )
                        Change volume.
                    else if ( $rand() \leq (nBox1 + nBox2 + nVol + 0.5nTrans)/nTotal$ )
                        Transfer particle from box 1 to box 2.
                    else
                        Transfer particle from box 2 to box 1.
                    end if
                end j loop
            end i loop

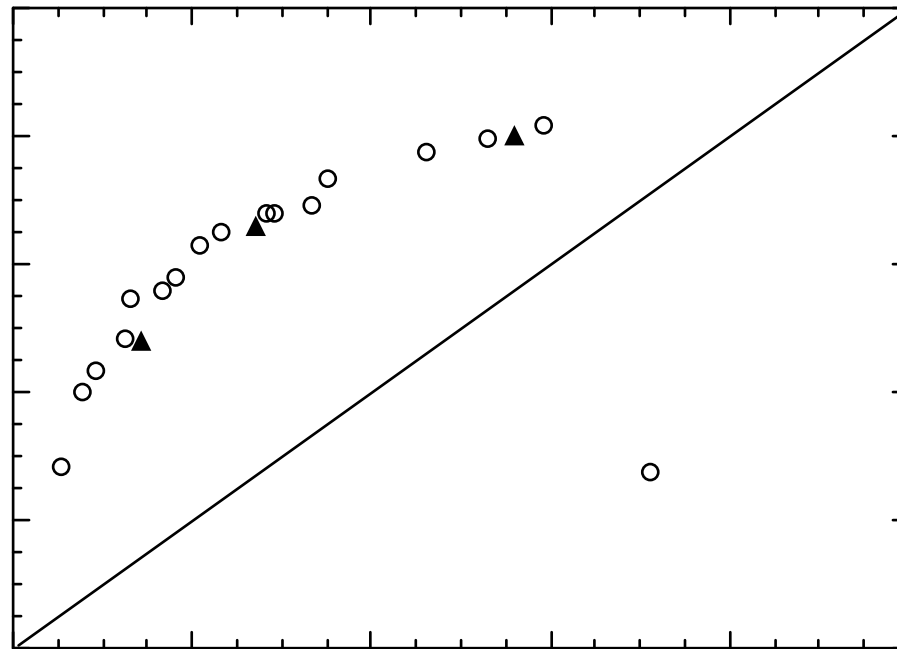
Part 2.1
Part 2.2
Part 2.3
```

Gibbs Ensemble – Result



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■ Water + Methanol Mixture



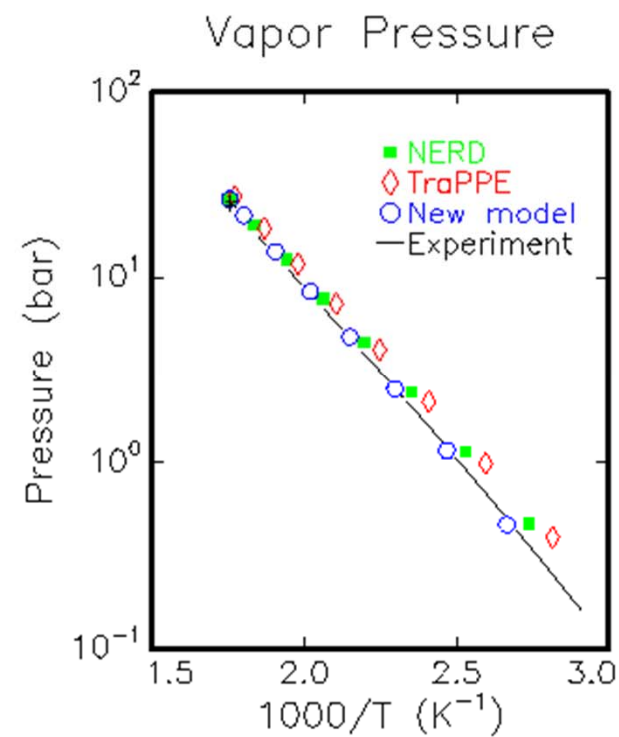
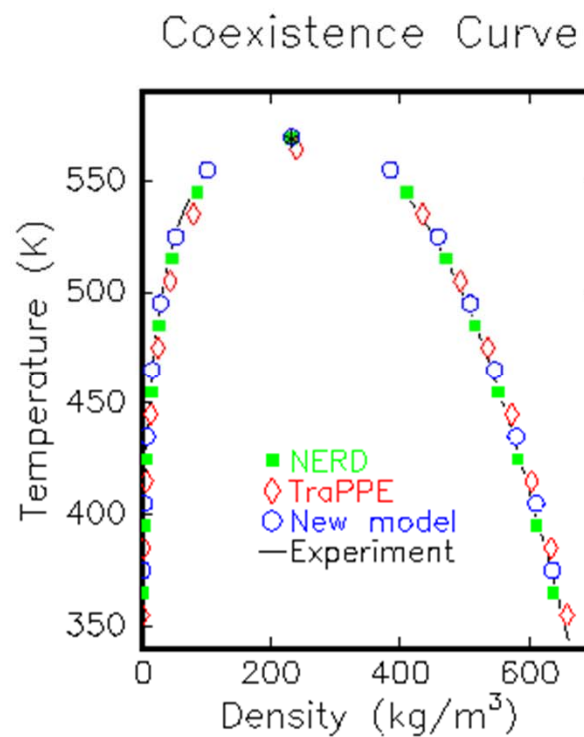
- *Strauch and Cummings, Fluid Phase Equilibria, 86 (1993) 147-172;*
- *Chialvo and Cummings, Molecular Simulation, 11 (1993) 163-175.*

Gibbs Ensemble – Result

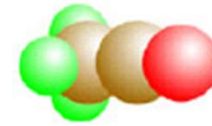


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■ Panagiotopoulos Group

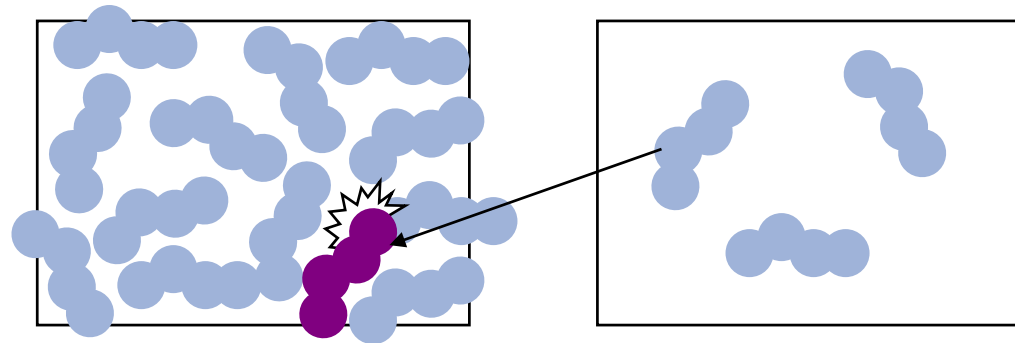


Gibbs Ensemble Limitation

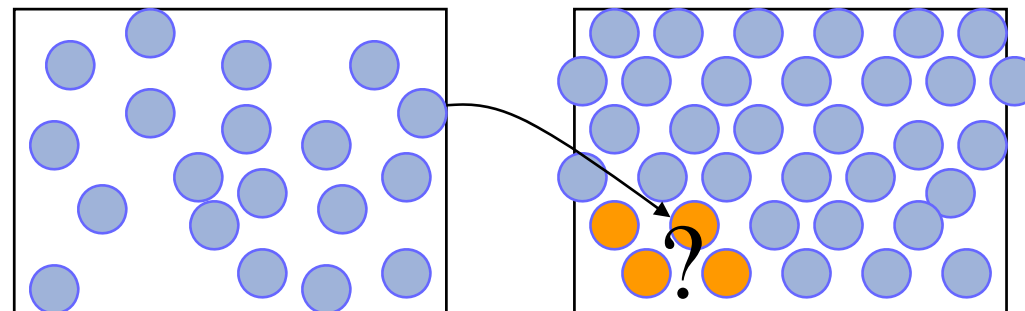


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- Limitation arise from particle-exchange requirements

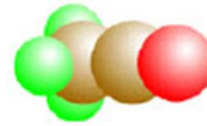


Dense phase or
Complex molecules



Solid phases

Gibbs-Duhem Integration



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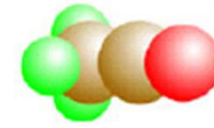
- David Kofke (1993)
- Basis : Numerical Integration of Cluisus – Clapeyron Equation

$$\mu_{\alpha} - \mu_{\beta} = -(S_{\alpha} - S_{\beta})dT - (V_{\alpha} - V_{\beta})dP$$

$$\frac{dP}{dT} = \frac{S_{\alpha} - S_{\beta}}{V_{\alpha} - V_{\beta}} = \frac{\Delta H}{T\Delta V}$$

$$\frac{d \ln P}{d1/T} = -\frac{\Delta H}{P\Delta V / T} \longrightarrow \left(\frac{\partial \ln p}{\partial \beta} \right)_{\sigma} = -\frac{\Delta h}{\Delta Z}$$

Gibbs – Duhem Integration

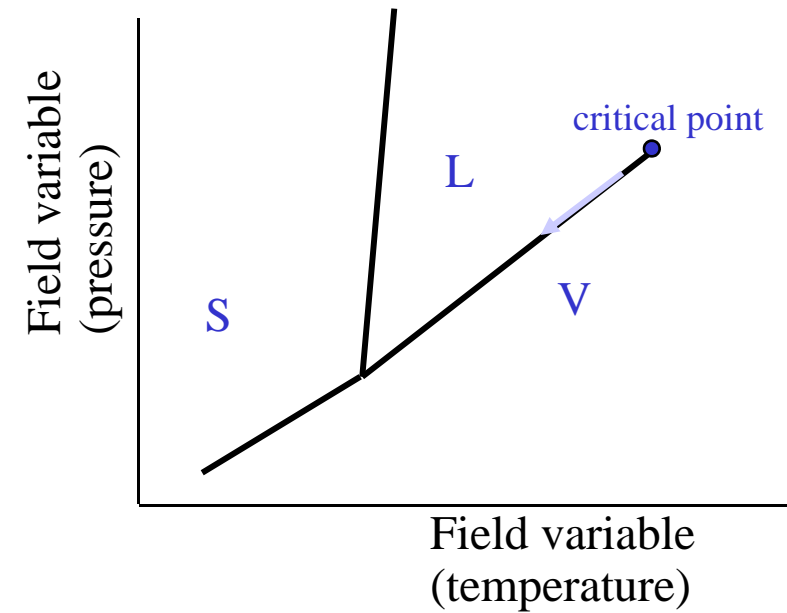


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- GE equation

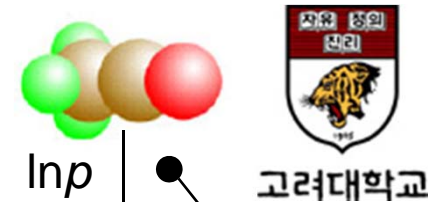
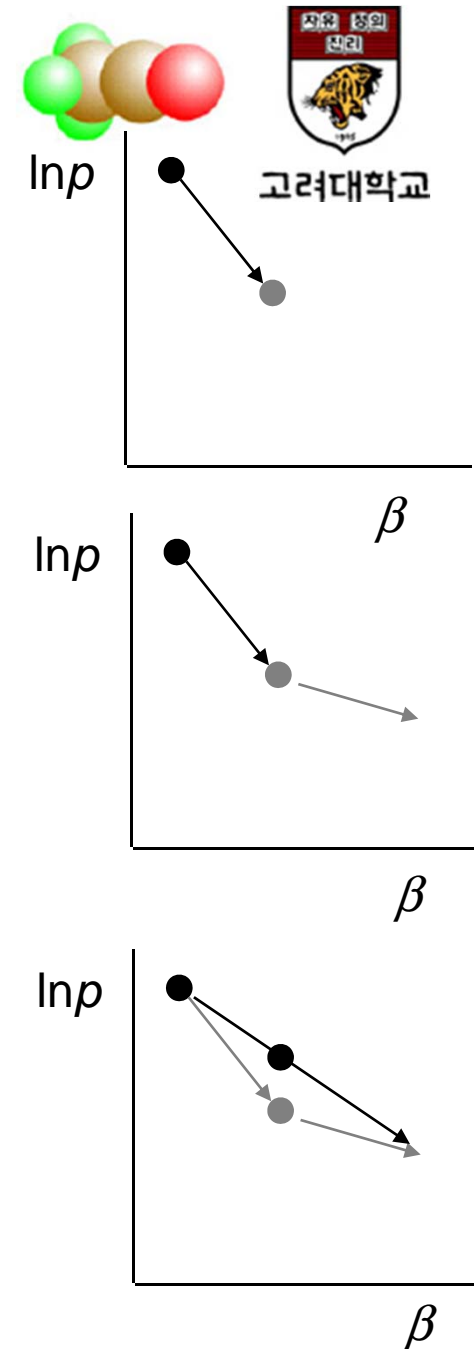
$$\left(\frac{\partial \ln p}{\partial \beta} \right)_\sigma = - \frac{\Delta h}{\Delta Z}$$

- Treat as nonlinear first order ODE
 - Use (NPT) simulation to obtain $\Delta H/\Delta z$

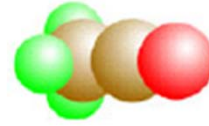


Predictor-Corrector Algorithm Implementation

- Given initial condition and slope ($= -\Delta h/\Delta Z$), predict new (p,T) pair.
- Evaluate slope at new state condition...
- ...and use to correct estimate of new (p,T) pair



Gibbs Duhem Integration



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- Potentially very efficient algorithm
 - Coexistence curves for solid-liquid systems
 - Coexistence curves for complex molecules
- Algorithm is not robust
- No built-in diagnostics

- Additional free energy calculation may be required to check the result