v For gases

Theory : $D_{\nu} = \frac{1}{3} \overline{u} \lambda$ $3 \frac{1}{2}$ $\bar{u}\lambda$ where, \bar{u} is average molar velocity, λ is mean free path Note, $\overline{u} \propto \frac{1}{R}$, $\therefore D_v \propto \frac{1}{R}$ $\frac{1}{p}$, \therefore $D_v \propto \frac{1}{p}$, $D_v \times P \sim$ $\frac{1}{p}$, $D_v \times P$ ~ constant (up to 1 atm) * \bar{u} $\propto T^{0.5}$, $\lambda \propto T^{1.5}$

Chapman-Enskog equation (for binary equation)

$$
D_{\nu} (= D_{AB}) \left[\frac{cm^2}{s} \right] = \frac{0.001858 T^{1.5} \left(\frac{M_A + M_B}{M_A M_B} \right)^{0.5}}{P[\alpha \text{tan} \cdot \sigma_{AB}[\hat{A}] \Omega_D} \qquad \Omega
$$

$$
\sigma_{AB} = \frac{\sigma_A + \sigma_B}{2} :
$$
 effective collision diameter
\n
$$
\Omega_D = f(kT/\epsilon_{AB}) :
$$
collision integral
\n
$$
\epsilon_{AB} = \sqrt{\epsilon_A \epsilon_B}
$$

온도상승에 따라 감소하는 값 하지만, 300~1,000K 에서 크게 변화하지 않음 \therefore D_{ν} 전체를 근사하면 \rightarrow D_{ν} ∝ $T^{1.75}$

Fuller equation

$$
D_{\nu} (= D_{AB}) \left[\frac{cm^2}{s} \right] = \frac{1.0110 \times 10^{-3} \cdot T^{1.5}}{P \left[(\Sigma V_A)^{\frac{1}{3}} + (\Sigma V_B)^{\frac{1}{3}} \right]^2} \left(\frac{1}{M_A} + \frac{1}{M_B} \right)^{0.5}
$$

 ΣV_i : Sum of diffusion volume of the component i from table

For pore size diffusion

• Knudsen diffusion : Diffusion in **VERY SMALL PORES** molecular collision on pore walls, the diffusivity \rightarrow less than normal volume

pore size $\ll \lambda$, thus pore size determines the diffusivity $D_K = 9{,}700r\sqrt{T/M}$ (for cylindrical pores)

• For intermediate-sized pores,

collisions with both pore walls & other molecules

$$
\frac{1}{D_{pore}} = \frac{1}{D_{AB}} + \frac{1}{D_K}
$$

❖ For liquids

Avg. travel dst. btw. collision is VERY LOW (→ 가 주로 분자크기보다 작음) therefore, D_{ν} in liquids are much smaller than those in gases.

 $D_n(l) \simeq 10^{-5} \sim 10^{-4} D_n(q)$

But Densities of : GASES << LIQUIDS ($\rightarrow \emptyset$ atmos. pressure), the fluxes for a given molar fraction gradient in liquid/gas may be nearly the same.

Stokes - Einstein equation : for Large & Spherical molecules in dilute solution

=> 유체의 흐름으로 인한 항력 (drag)' 고려 => the simplest equation

$$
D_{\nu} = \frac{kT}{6\pi r_{0}\mu} \simeq 7.32 \times 10^{-16} \frac{T}{r_{0}\mu} \qquad D_{\nu} \propto \frac{1}{V^{1/3}}
$$

.

2) Prediction of diffusivity

Wilke-Chang equation : for solutes of small to moderate size, => the $D_{\nu}(l)$ becomes greater as the drag is less than predicted

 \triangleright the empirical equation is:

$$
D_{\nu} = 7.4 \times 10^{-8} \frac{T \sqrt{\psi_B M_B}}{V_A^{0.6} \mu} \qquad D_{\nu} \propto \frac{1}{V^{1/6}}
$$

A: solute, B: solvent

 ψ_B : association parameter (to be given)

17.2 Turbulent diffusion

Turbulent \rightarrow Eddies transport matter ! (transfer momentum & heat energy)

17.3 Mass transfer coefficient

 δ (film thickness) Film theory

17.3 Mass transfer coefficient

① equimolar diffusion

•
$$
N_A = k_c \Delta C_A = k_y \Delta y_A
$$

•
$$
N_A = D_v \frac{\Delta C_A}{\delta} = D_v C \frac{\Delta y_A}{\delta}
$$

•
$$
k_c = \frac{D_v}{\delta}, k_y = \frac{D_v}{\delta} \cdot C (= k_c \cdot C)
$$

② one-waydiffusion

•
$$
N_A = k_c' \Delta C_A = k_y' \Delta y_A
$$

\n• $N_A = D_v \frac{C}{(C - C_A)m} \frac{\Delta C_A}{\delta} \Rightarrow D_v \frac{1}{(1 - y_A)m} \frac{\Delta C_A}{\delta}$
\n $= D_v C \frac{1}{(1 - y_A)m} \frac{\Delta y_A}{\delta}$
\n• $k_c' = \frac{D_v}{\delta} \frac{1}{(1 - y_A)m}, \dots, k_c' = k_c \frac{1}{(1 - y_A)m}$
\n $k_y' = \frac{D_v \cdot C}{\delta} \frac{1}{(1 - y_A)m}, \dots, k_y' = k_y \frac{1}{(1 - y_A)m}$
\n $\therefore \frac{k_y'}{k_y} = \frac{k_c'}{k_c} \frac{C}{C} \left(= \frac{1}{(1 - y_A)m} \right)$

